

# Methods for optical calibration of the BigBite hadron spectrometer

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## Abstract

The techniques for optical calibration of Jefferson Lab's large-acceptance magnetic hadron spectrometer, BigBite, have been examined. The most consistent and stable results were obtained by using a method based on singular value decomposition. In spite of the complexity of the optics, the particles' positions and momenta at the target have been precisely reconstructed from the coordinates measured in the detectors by means of a single back-tracing matrix. The technique is applicable to any similar magnetic spectrometer and any particle type. For 0.55 GeV/c protons, we have established a vertex resolution of 1.2 cm, angular resolutions of 7 mrad and 13 mrad (in-plane and out-of-plane, respectively), and a relative momentum resolution of 1.6 %.

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## 1. Introduction

One of the recent acquisitions in experimental Hall A of the Thomas Jefferson National Accelerator Facility (TJNAF) is the BigBite spectrometer. It was previously used at the NIKHEF facility for the detection of electrons [1, 2]. At Jefferson Lab, BigBite has been re-implemented as a versatile spectrometer that can be instrumented with various detector packages optimized for the particular requirements of the experiments. BigBite complements the High-Resolution Spectrometers, which are part of the standard equipment of Hall A [3]. Adding BigBite allows one to devise more flexible experimental setups involving double- and even triple-coincidence measurements.

In 2005, the BigBite spectrometer was first used in Hall A as the hadron arm in the E01-015 experiment, which investigated nucleon-nucleon short-range correlations [4, 5]. In 2006, it was instrumented as the electron arm for the measurement of the neutron electric form factor (experiment E02-013 [6]). In 2008 and 2009, it has been used in two large groups of experiments spanning a broad range of physics topics. We studied near-threshold neutral pion production on protons (experiment E04-007 [7]) and measured single-spin asymmetries in semi-inclusive pion electro-production on polarized  $^3\text{He}$  (experiments E06-010 and E06-011 [8, 9]). In the same period, we also measured parallel and perpendicular asymmetries on polarized  $^3\text{He}$  in order to extract the  $g_2^n$  polarized structure function in the deep-inelastic regime (experiment E06-014 [10]), and measured double-polarization asymmetries in the quasi-elastic processes  $^3\vec{\text{He}}(\vec{e}, e'd)$ ,  $^3\vec{\text{He}}(\vec{e}, e'p)$ , and  $^3\vec{\text{He}}(\vec{e}, e'n)$  (experiments E05-102 and E08-005 [11, 12]). In 2011, the investigation of short-range correlations has been continued in the E07-006 experiment [13] exploring the repulsive part of the nucleon-nucleon interaction.

BigBite is a non-focusing spectrometer consisting of a single dipole with large momentum and angular acceptances (the details are presented in Section 2). The magnetic optics of such spectrometers tend to become complicated towards the edges of their acceptances, especially for the momentum and the dispersive angle. It was not clear from the outset that particle momentum and interaction vertex reconstruction could be accomplished by using a single procedure for all momenta.

The calibration presented in this paper allows for a full description of BigBite optics by means of a single reconstruction matrix. The method was developed and successfully used

with the data obtained in the E05-102 experiment with the detector package configured for hadrons (Section 3), but it is applicable to any magnetic spectrometer with a similar optical configuration and any particle type. Various calibration procedures are discussed in Section 4.

## 2. The BigBite spectrometer

The BigBite spectrometer [1] consists of a single room-temperature dipole magnet, shown in Fig. 1. Energizing the magnet with a current of 518 A results in a mean field density of 0.92 T, corresponding to a central momentum of  $p_c = 0.5 \text{ GeV}/c$  and a bending angle of 25°. The magnet is coupled to a hadron detector package consisting of two multi-wire drift chambers (MWDC) [14, 15] for particle tracking and two planes of scintillation detectors (denoted by dE and E) [16] for triggering, particle identification, and energy determination.

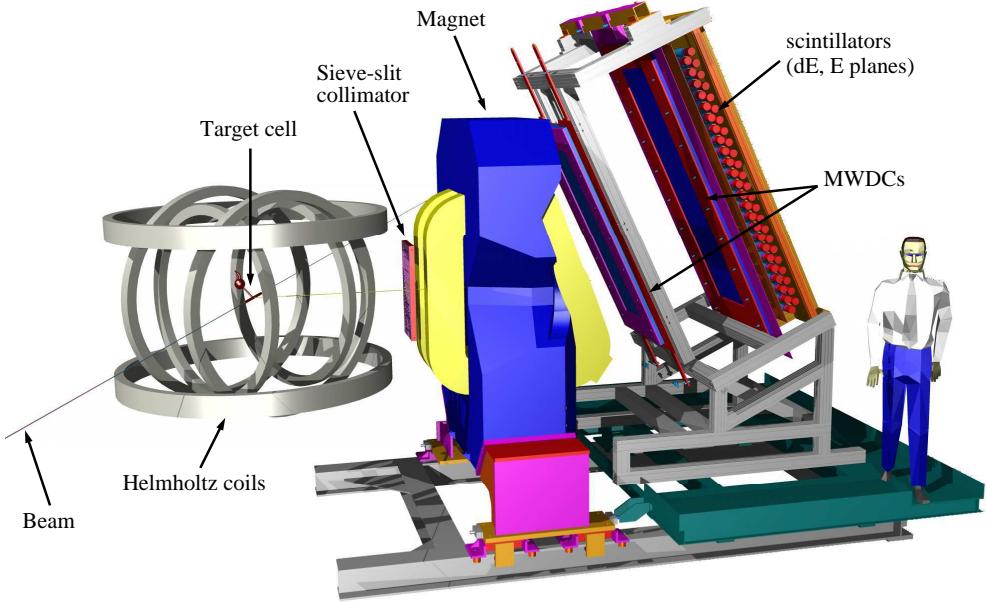


Figure 1: The BigBite spectrometer on its support frame. BigBite consists of a dipole magnet, followed by the detector package assembled from a pair of multi-wire drift chambers (MWDC) and two scintillator planes (dE and E). The directions of the incoming electron beam and the scattered particles, the target cell, and the Helmholtz coil (holding field) assembly are also shown.

Each MWDC consists of six planes of wires. The wires in the first two planes are oriented at an angle of  $60^\circ$  with respect to the dispersive direction, while the wires in the third and fourth plane are aligned horizontally. The wires of the last two planes are oriented at  $-60^\circ$ . Each wire plane in the first and the second MWDC contains 141 and 200 wires, respectively. The spacing between the wires in all planes is 1 cm. The intrinsic spatial resolution of the MWDCs is about  $100\ \mu\text{m}$  and  $200\ \mu\text{m}$  for the dispersive and non-dispersive coordinates, respectively, and about 0.15 mrad and 0.35 mrad for the dispersive and non-dispersive angles, respectively.

The dE- and E-planes (also called the trigger planes) each consist of 24 plastic scintillator bars. The bars are 50 cm long and 8.6 cm wide. For the dE-plane, thinner bars (0.3 cm) were used to detect low-energy particles, while for the E-plane, a thickness of 3 cm was chosen to allow for the detection of more energetic particles. The light pulses in each bar were detected by photomultiplier tubes mounted at each end of the bar. To double the spatial and momentum resolution, the bars in the E-plane are offset from those in the dE-plane by one half of the bar width (4.3 cm).

### 3. Experimental details and data

The E05-102 experiment was performed in Hall A [3] at Jefferson Lab. In the experiment, a polarized  $^3\text{He}$  target was used in conjunction with the polarized continuous-wave electron beam. Scattered electrons were detected by the left High Resolution Spectrometer (HRS) in coincidence with protons and deuterons that were detected by BigBite. A variety of kinematic settings were employed (Table 1), with the momentum-transfer vector  $\vec{q}$  pointing towards BigBite. This ensured that the protons and deuterons from elastic and quasi-elastic scattering were always within its acceptance.

The core component of the polarized  $^3\text{He}$  target was a pressurized cylindrical glass cell with a length of 40 cm and a diameter of 1.9 cm (see Fig. 2). The thickness of the glass cylinder was 1.7 mm, while the thickness of the end windows was  $140\ \mu\text{m}$ . The gas in the cell was polarized to approximately 60 % by hybrid spin-exchange optical pumping [17, 18] driven by an infra-red laser system. The direction of the nuclear polarization was maintained by three pairs of Helmholtz coils surrounding the cell.

Table 1: Kinematics settings of the E05-102 experiment (the incoming electron energy  $E_{\text{beam}}$  and the angles of the HRS and BigBite spectrometers with respect to the beam direction).

Setting label	$E_{\text{beam}}$ [GeV/c]	Scattering angle HRS [°]	BB [°]
1-pass	1.245	17.0	-74.0
2-pass	2.425	12.5	-75.0
		14.5	-82.0
3-pass	3.606	12.5	-75.0
		17.0	-74.0

In addition to the  ${}^3\text{He}$  helium target, a 40 cm-long multi-foil carbon target was used for calibration, as described below. It consists of seven 0.252 mm-thick carbon foils mounted to a plastic frame (Fig. 2) which are preceded by a single slanted BeO foil for beam positioning. Below the multi-foil target, a dummy (reference) cell was installed that could be either evacuated or filled with hydrogen, deuterium, unpolarized helium-3 or nitrogen.

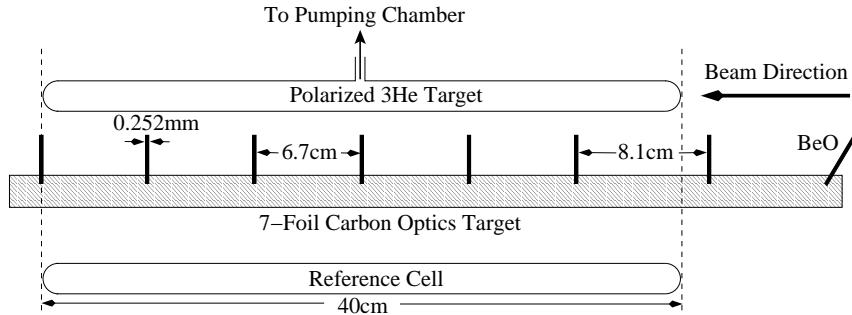


Figure 2: The target system including the polarized  ${}^3\text{He}$  cell at the top, the multi-foil carbon optics target, and the reference cell at the bottom. The slanted BeO foil is used for visual inspection of the beam impact point.

For the optics calibration of BigBite, a special set of measurements was performed with a 4 cm-thick lead sieve-slit collimator positioned at the entrance to the BigBite magnet (see Fig. 1). The sieve-slit collimator has 82 circular holes that are almost uniformly positioned over the whole acceptance of the spectrometer, Fig. 3 (left). The

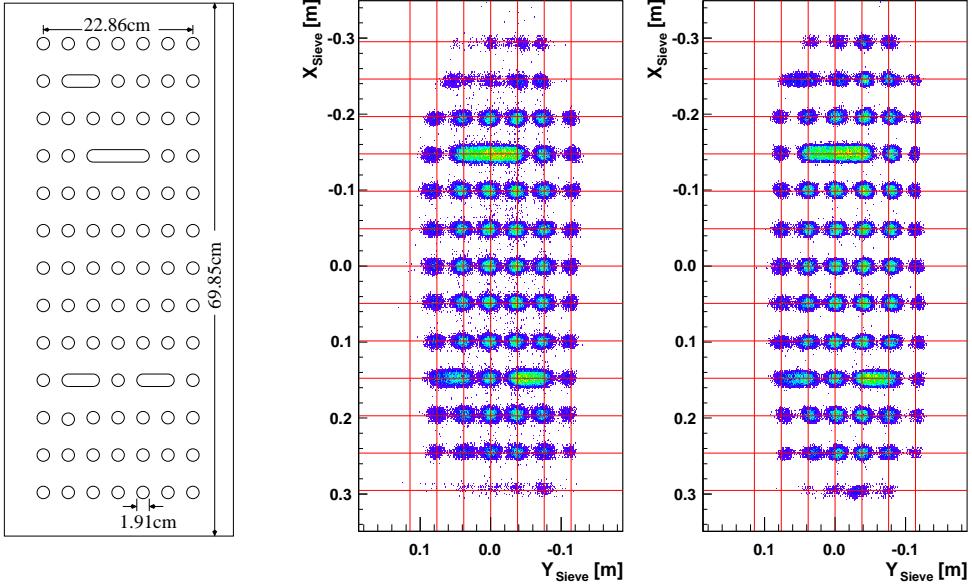


Figure 3: [Left] Schematics of the BigBite sieve-slit collimator. [Center, Right] Sieve pattern reconstruction by using the simplex method and the SVD, respectively (see subsubsection 4.2.2). The SVD technique resolves more holes and yields a much clearer pattern. The holes at the left edge are missing due to geometrical obstacles between the target and BigBite.

collimator also contains four elongated holes used to remove ambiguities in horizontal and vertical orientations and to allow for easier identification of the hole projections at the detector package.

Prior to any optics analysis, a series of cuts were applied to the collected calibration data to eliminate backgrounds. A HRS-BigBite coincidence trigger system was used to acquire electron-proton and electron-deuteron coincidences, at typical rates between 700 Hz and 1 kHz. True coincidences were selected by applying a cut on the raw coincidence time. The backgrounds were further reduced by PID and HRS acceptance cuts. Finally, only those events that produce consistent hits in all BigBite detectors, and could consequently be joined to form single particle tracks, were selected.

## 4. Methods of optical calibration

The purpose of optical calibration is to establish the mapping between the detector variables that are measured directly, and the target variables corresponding to the actual physical quantities describing the particle at the reaction vertex. In the MWDCs, two position coordinates ( $x_{\text{Det}}$  and  $y_{\text{Det}}$ ) and two angles ( $\theta_{\text{Det}}$  and  $\phi_{\text{Det}}$ ) are measured. From this information, we wish to reconstruct the location of the interaction vertex ( $y_{\text{Tg}}$ ), the in-plane and out-of-plane scattering angles ( $\phi_{\text{Tg}}$  and  $\theta_{\text{Tg}}$ ), and the particle momentum relative to the central momentum ( $\delta_{\text{Tg}} = (p_{\text{Tg}} - p_c)/p_c$ ). This can be done in many ways. We have considered an analytical model as well as a more sophisticated approach based on transport-matrix formalism, with several means to estimate the reliability of the results and the stability of the algorithms.

Quasi-elastic protons from scattering on the multi-foil carbon target were used to calibrate  $y_{\text{Tg}}$ ; the same target was also used to calibrate  $\theta_{\text{Tg}}$  and  $\phi_{\text{Tg}}$  when the sieve-slit collimator was in place. In turn, elastic protons and deuterons (from hydrogen and deuterium targets) were used to calibrate  $\theta_{\text{Tg}}$ ,  $\phi_{\text{Tg}}$ , and  $\delta_{\text{Tg}}$ . The  $\delta_{\text{Tg}}$  matrix elements could also be determined by quasi-elastic events from  ${}^3\text{He}$  under the assumption that the energy losses are well understood.

### 4.1. The analytical model

The magnetic field of the BigBite magnet is oriented in the  $y_{\text{Tg}}$  direction (see Fig. 4). Field mapping has shown [1] that the field density is almost constant inside the magnet, with fringe fields that decrease exponentially outside of the magnet. In the analytical model, the true field was approximated by a constant field within the effective field boundaries, while edge effects were neglected. Under these assumptions all target coordinates were calculated by applying a circular-arc approximation [21] of the track inside the field: the particle transport was divided into free motion (drift) in the  $(y, z)$  plane and circular motion in the  $(x, z)$  plane (see Fig. 4), described by the Lorentz equation

$$p_y = \text{const} , \quad p_{xz} = eRB_y .$$

To determine the momentum, the radius  $R$  of the trajectory needs to be calculated first. This can be done by using the track information obtained from the detector package,

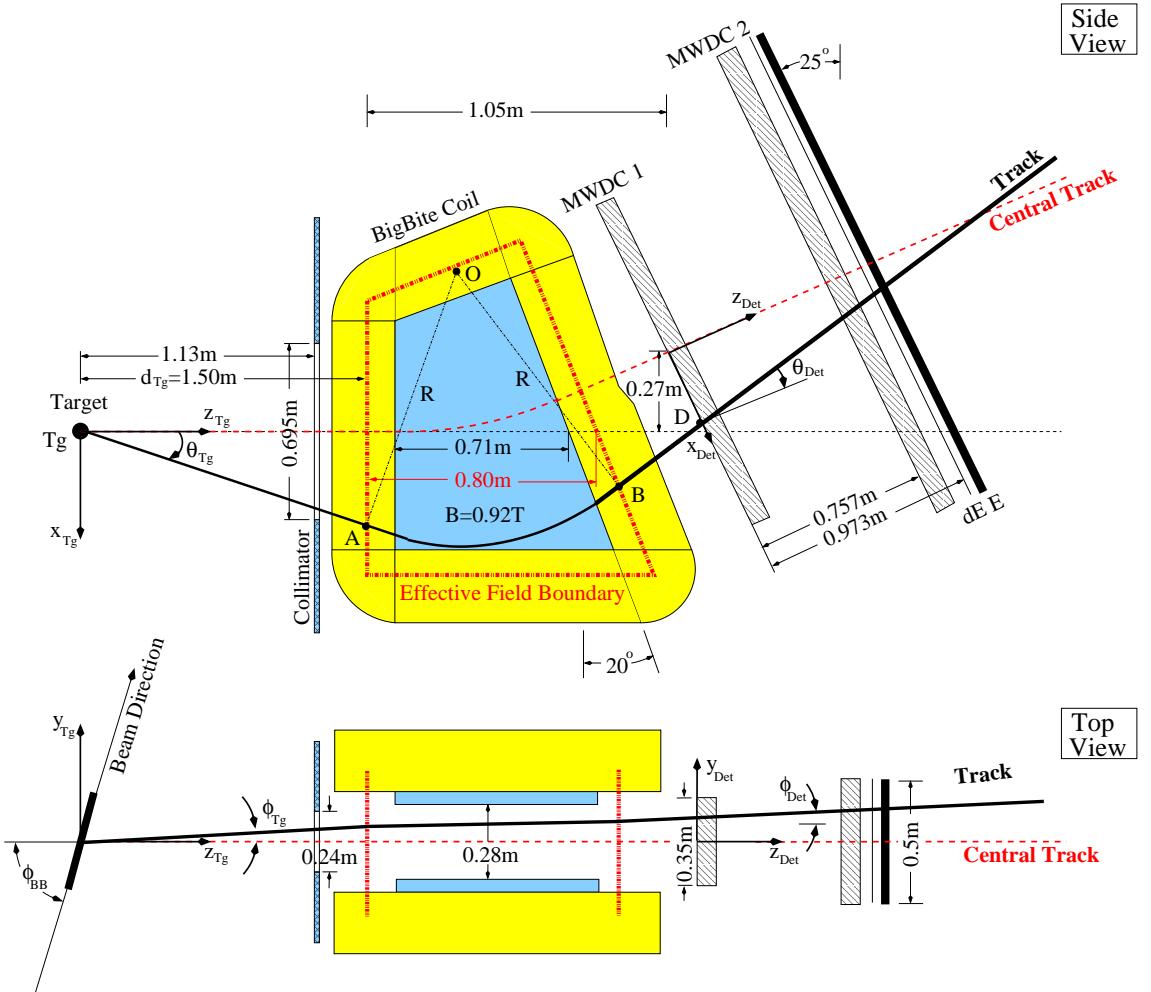


Figure 4: The schematic of the dispersive (top) and non-dispersive (bottom) planes of the BigBite spectrometer. Small angular deflections in the non-dispersive plane occur if the particle trajectory is not perpendicular to the effective field boundary [2, 19, 20]. At the entrance to the magnet, they are at most 18 mrad (close to the acceptance boundaries in the dispersive direction). At the exit field boundary, the effect acts in the opposite sense and partially cancels the deflection at the entrance.

combined with the geometrical properties of BigBite. A few reference points are needed, as shown in Fig. 4; the point  $T_g$  represents the position of the particle at the target, and  $D$  corresponds to the point where the particle hits the detector package. The point  $B$  at which the particle exits the magnet is the intersection between the extrapolated particle track through the detector package and the effective exit face of the magnet. Similarly, the point  $A$  lies at the intersection of the effective entrance face of the magnet and the particle track from the target. The point  $O$  is the center of the circular trajectory. In order for all these points to correspond to a single particle track through the spectrometer, the conditions

$$\overline{AT_g} \perp \overline{AO}, \quad \overline{OB} \perp \overline{BD}, \quad |\overline{AO}| = |\overline{BO}| = R,$$

must be satisfied. In the target coordinate system, this becomes

$$x_O = -\frac{d_{Tg}}{(x_A - x_{Tg})}(z_O - z_A) + x_A = -\frac{z_D - z_B}{x_D - x_B}(z_O - z_B) + x_B, \quad (1)$$

$$R^2 = (z_O - z_A)^2 \left[ 1 + \left( \frac{d_{Tg}}{x_A - x_{Tg}} \right)^2 \right] = (z_B - z_O)^2 \left[ 1 + \left( \frac{z_D - z_B}{x_D - x_B} \right)^2 \right]. \quad (2)$$

The coordinates  $x_B$  and  $z_B$  of  $B$ , and the coordinates  $x_D$  and  $z_D$  of  $D$  can be directly calculated from the information obtained by the detector package. The position of the target  $(x_{Tg}, z_{Tg})$  is known. Since only thin targets are employed,  $x_{Tg}$  is set to zero. The coordinate  $z_A$  of  $A$  corresponds to the known distance  $d_{Tg}$  between the target center and the effective field boundary at the entrance to the magnet. By expressing  $z_O$  from Eq. (1) and inserting it into Eq. (2), an equation for  $x_A$  is obtained which has three complex solutions in general. The physically meaningful result for  $x_A$  should be real and lie within the effective field boundaries. Two additional physical constraints are applied. The particle track should always represent the shortest possible arc of the circle (the arc between  $A$  and  $B$  in Fig. 4). Moreover, the track should bend according to the polarity of the particle and orientation of the magnetic field. From  $x_A$ , the radius  $R$  and the momentum  $p_{xz}$  can be calculated. The particle flight path  $l_{xz}$  in the  $(x, z)$  plane can also be calculated by using the cosine formula for the angle  $\beta = \angle AOB$ ,

$$l_{xz} = \sqrt{x_A^2 + d_{Tg}^2} + R\beta + \sqrt{(x_D - x_B)^2 + (z_D - z_B)^2},$$

$$\cos \beta = \frac{(x_A - x_O)(x_B - x_O) + (z_A - z_O)(z_B - z_O)}{R^2}.$$

By using this information, all target coordinates can be expressed as

$$\begin{aligned}
\phi_{\text{Tg}} &= \phi_{\text{Det}}, \\
\theta_{\text{Tg}} &= \arctan\left(\frac{x_A}{d_{\text{Tg}}}\right), \\
y_{\text{Tg}} &= y_{\text{Det}} - l_{xz} \tan \phi_{\text{Det}}, \\
\delta_{\text{Tg}} &= \frac{p_{xz}}{p_c} \frac{\sqrt{1 + \tan^2 \phi_{\text{Tg}} + \tan^2 \theta_{\text{Tg}}}}{\sqrt{1 + \tan^2 \theta_{\text{Tg}}}} - 1, \\
L &= l_{xz} \sqrt{1 + \tan^2 \phi_{\text{Tg}}},
\end{aligned}$$

where  $p_c$  is the central momentum and  $L$  is the total flight-path of the particle.

With the analytical approximation, resolutions of a few percent can be achieved, but they deteriorate when moving towards the edges of the acceptance where the fringe fields begin to affect the optics. This is particularly true for  $\phi_{\text{Tg}}$ . Figure 5 (left) shows the reconstructed mass of the neutron from the process  $^2\text{H}(\text{e}, \text{e}'\text{p})\text{n}$ , obtained by using the analytical model. The relative resolution is 0.35 %.

The analytical method requires just a few geometry parameters, but these need to be known quite accurately. Had no survey been performed, the sizes of spectrometer components and the distances between them could be obtained, in principle, by calibrating with elastic events. However, the solution is not unique. Different combinations of parameters have been shown to yield almost identical results for the target variables, while only one combination is correct.

#### 4.2. The matrix formalism

In spite of its shortcomings, the analytical model is a good starting point. Due to its simplicity, it can be implemented and tested quickly, and lends itself well to online estimation of the experimental data. For the off-line analysis, a more sophisticated approach based on the transport matrix formalism is needed. In this approach, a prescription is obtained that transforms the detector variables directly to the target variables. Various parameterizations of this transformation are possible. We have adopted a polynomial expansion of the form [22, 23]

$$\Omega_{\text{Tg}} = \sum_{i,j,k,l} a_{ijkl}^{\Omega_{\text{Tg}}} x_{\text{Det}}^i \theta_{\text{Det}}^j y_{\text{Det}}^k \phi_{\text{Det}}^l, \quad \Omega_{\text{Tg}} \in \{\delta_{\text{Tg}}, \theta_{\text{Tg}}, \phi_{\text{Tg}}, y_{\text{Tg}}\}. \quad (3)$$

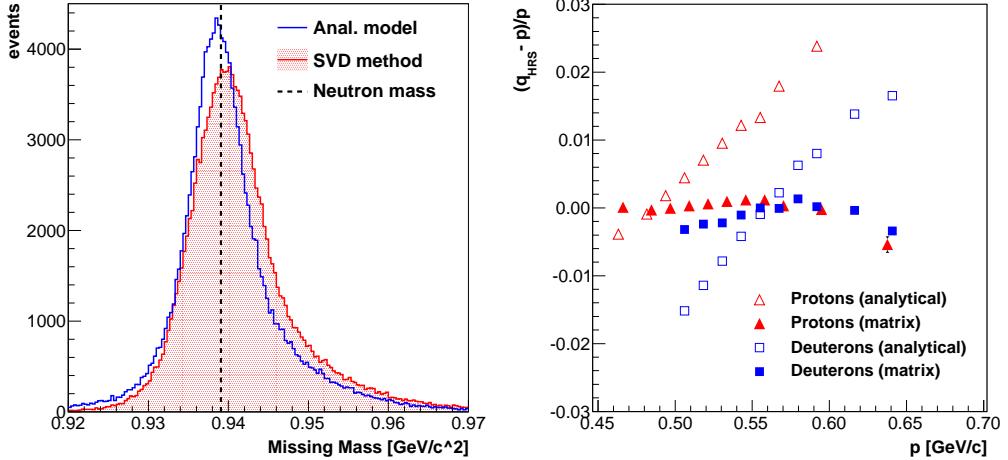


Figure 5: [Left] The reconstructed mass of the undetected neutron (missing mass) from the process  $^2\text{H}(\text{e}, \text{e}'\text{p})\text{n}$  by using the analytical model and the matrix-formalism (SVD) approach. The width (sigma) of the peak determined with the analytical model is  $3.3 \text{ MeV}/c^2$  (corresponding to 0.35 % relative resolution). The width of the peak reconstructed by the SVD method is  $4 \text{ MeV}/c^2$ . [Right] The absolute calibration of  $\delta_{Tg}$  as a function of the particle momentum measured by BigBite. The relative resolution of  $\delta_{Tg}$  is better in the analytical model than in the matrix method, but the absolute momentum calibration is inferior to the matrix approach, except in the narrow region around  $p \approx 0.55 \text{ GeV}/c$ .

Knowing the optics of a spectrometer is equivalent to determining the expansion coefficients  $a_{ijkl}^{\Omega_{Tg}}$  (the so-called optical ‘‘matrix’’) and establishing the limitations of such a parameterization.

Ideally, one would like to obtain a single optical matrix with full reconstruction functionality for all particle species and momenta, with as few high-order terms as possible. In a large-acceptance spectrometer like BigBite, this represents a considerable challenge. In particular, one must clearly understand the contributions of the high-order elements. Uncontrolled inclusion of these terms typically causes oscillations of the reconstructed variables at the edges of the acceptance. In the following sections, we discuss the procedure of constructing the optical matrix in which special attention is devoted to checking the convergence of the method and estimating the robustness of the matrix elements.

#### 4.2.1. Decoupled description

The determination of the optical matrix starts with a low-order analysis in order to estimate the dominant matrix elements. As in the analytical model, the BigBite magnet is assumed to be an ideal dipole. This assumption decouples the in-plane and out-of-plane variables, resulting in the simplification that  $\delta_{\text{Tg}}$  and  $\theta_{\text{Tg}}$  depend only on  $x_{\text{Det}}$  and  $\theta_{\text{Det}}$ , while  $y_{\text{Tg}}$  and  $\phi_{\text{Tg}}$  depend only on  $y_{\text{Det}}$  and  $\phi_{\text{Det}}$ .

Since each target coordinate depends only on two detector coordinates, the matrix elements were estimated by examining two-dimensional histograms of target coordinates (as given by the HRS) versus BigBite detector variables, using various detector-variable cuts. Since BigBite in this approximation does not bend horizontally, only first-order polynomials were utilized to fit the data for  $y_{\text{Tg}}$  and  $\phi_{\text{Tg}}$ , while expansions up to third-order were applied for  $\delta_{\text{Tg}}$  and  $\theta_{\text{Tg}}$ :

$$\begin{aligned}\delta_{\text{Tg}}(x, \theta) &= \left[ a_{0000}^{\delta_{\text{Tg}}} + a_{1000}^{\delta_{\text{Tg}}}x + a_{2000}^{\delta_{\text{Tg}}}x^2 \right] + \left[ a_{0100}^{\delta_{\text{Tg}}} + a_{1100}^{\delta_{\text{Tg}}}x + a_{2100}^{\delta_{\text{Tg}}}x^2 \right] \theta \\ &+ \left[ a_{0200}^{\delta_{\text{Tg}}} + a_{1200}^{\delta_{\text{Tg}}}x \right] \theta^2 + \left[ a_{0300}^{\delta_{\text{Tg}}} + a_{1300}^{\delta_{\text{Tg}}}x \right] \theta^3, \\ \theta_{\text{Tg}}(x, \theta) &= \left[ a_{0000}^{\theta_{\text{Tg}}} + a_{1000}^{\theta_{\text{Tg}}}x + a_{2000}^{\theta_{\text{Tg}}}x^2 \right] + \left[ a_{0100}^{\theta_{\text{Tg}}} + a_{1100}^{\theta_{\text{Tg}}}x + a_{2100}^{\theta_{\text{Tg}}}x^2 \right] \theta, \\ \phi_{\text{Tg}}(y, \phi) &= a_{0000}^{\phi_{\text{Tg}}} + a_{0001}^{\phi_{\text{Tg}}}\phi, \\ y_{\text{Tg}}(y, \phi) &= \left[ a_{0001}^{y_{\text{Tg}}} + a_{0011}^{y_{\text{Tg}}}y \right] \phi + \left[ a_{0000}^{y_{\text{Tg}}} + a_{0010}^{y_{\text{Tg}}}y \right].\end{aligned}$$

The calculated matrix elements are shown in the second column of Table 2. The  $a_{0001}^{\phi_{\text{Tg}}}$  matrix element was set to 1 since there is no in-plane bending. This approximation could not be used for further physics analysis because higher-order corrections are needed. However, the low-order terms are very robust and do not change much when more sophisticated models with higher-order terms are considered. The results obtained by using this method serve as a benchmark for more advanced methods, in particular as a check whether the matrix elements computed by automated numerical algorithms converge to reasonable values.

#### 4.2.2. Higher order matrix formalism

For the determination of the optics matrix a numerical method was developed in which matrix elements up to fourth order were retained. Their values were calculated by using a  $\chi^2$ -minimization scheme, wherein the target variables calculated by Eq. (3) were

Table 2: The dominant matrix elements of the BigBite optics model (Eq. (3)) determined by a decoupled description (subsubsection 4.2.1), by simplex minimization (N&M), and by singular value decomposition (SVD, subsubsection 4.2.2).

Matrix element	Decoupled description	N&M	SVD
$a_{0010}^{y_{\text{Tg}}}$ [m/m]	0.998	1.024	0.917
$a_{0001}^{y_{\text{Tg}}}$ [m/rad]	-2.801	-2.839	-2.766
$a_{0001}^{\phi_{\text{Tg}}}$ [rad/rad]	1.000	1.052	0.9517
$a_{1000}^{\theta_{\text{Tg}}}$ [rad/m]	0.497	0.549	0.551
$a_{0100}^{\theta_{\text{Tg}}}$ [rad/rad]	-0.491	-0.490	-0.484
$a_{1000}^{\delta_{\text{Tg}}}$ [1/m]	-0.754	-0.716	-0.676
$a_{0100}^{\delta_{\text{Tg}}}$ [1/rad]	2.811	2.881	2.802

compared to the directly measured values,

$$\chi^2 \left( a_i^{\Omega_{\text{Tg}}} \right) = \sqrt{\left( \Omega_{\text{Tg}}^{\text{Measured}} - \Omega_{\text{Tg}}^{\text{Optics}} \left( x_{\text{Det}}, y_{\text{Det}}, \theta_{\text{Det}}, \phi_{\text{Det}}, a_i^{\Omega_{\text{Tg}}} \right) \right)^2}, \quad i = 1, 2, \dots, M. \quad (4)$$

The use of  $M$  matrix elements for each target variable means that a global minimum in  $M$ -dimensional space must be found. Numerically this is a very complex problem; two techniques were considered for its solution.

Our first choice was the downhill simplex method developed by Nelder and Mead [24, 25]. The method tries to minimize a scalar non-linear function of  $M$  parameters by using only function evaluations (no derivatives). It is widely used for non-linear unconstrained optimization, but it is inefficient and its convergence properties are poorly understood, especially in multi-dimensional minimizations. The method may stop in one of the local minima instead of the global minimum [27, 28], so an additional examination of the robustness of the method was required.

The set of functions  $\Omega_{\text{Tg}}$  is linear in the parameters  $a_i^{\Omega_{\text{Tg}}}$ . Therefore, Eq. (4) can be written as

$$\chi^2 = \sqrt{\left| A \vec{a} - \vec{b} \right|^2}, \quad (5)$$

where the  $M$ -dimensional vector  $\vec{a}$  contains the matrix elements  $a_i^{\Omega_{\text{Tg}}}$ , and the  $N$ -

dimensional vector  $\vec{b}$  contains the measured values of the target variable being considered. The elements of the  $N \times M$  matrix  $A$  are various products of detector variables ( $x_{\text{Det}}^i \theta_{\text{Det}}^j y_{\text{Det}}^k \phi_{\text{Det}}^l$ ) for each measured event. The system  $A \vec{a} = \vec{b}$  in Eq. (5) is overdetermined ( $N > M$ ), thus the vector  $\vec{a}$  that minimizes the  $\chi^2$  can be computed by singular value decomposition (SVD). It is given by  $A = UWV^T$ , where  $U$  is a  $N \times M$  column-orthogonal matrix,  $W$  is a  $M \times M$  diagonal matrix with non-negative singular values  $w_i$  on its diagonal, and  $V$  is a  $M \times M$  orthogonal matrix [25, 26]. The solution has the form

$$\vec{a} = \sum_{i=1}^M \left( \frac{\vec{U}_i \cdot \vec{b}}{w_i} \right) \vec{V}_i.$$

The SVD was adopted as an alternative to simplex minimization since it produces the best solution in the least-square sense, obviating the need for robustness tests. Another great advantage of SVD is that it can not fail; the method always returns a solution, but its meaningfulness depends on the quality of the input data. The most important leading-order matrix elements computed by using both techniques are compared in Table 2.

## 5. Calibration results

### 5.1. Vertex position

The matrix for the vertex position variable  $y_{\text{Tg}}$  was obtained by analyzing the protons from quasi-elastic scattering of electrons on the multi-foil carbon target. The positions of the foils were measured by a geodetic survey to sub-millimeter accuracy, allowing for a very precise calibration of  $y_{\text{Tg}}$ . The vertex information from the HRS was used to locate the foil in which the particle detected by BigBite originated. This allowed us to directly correlate the detector variables for each coincidence event to the interaction vertex. When Eq. (3) is written for  $y_{\text{Tg}}$ , a linear equation for each event can be formed:

$$\begin{aligned} y_{\text{Tg}(n)}^{\text{Measured}} = y_{\text{Tg}(n)}^{\text{Optics}} &:= a_{0000}^y + a_{0001}^y \phi_{(n)} + a_{0002}^y \phi_{(n)}^2 + a_{0003}^y \phi_{(n)}^3 + \dots \\ &+ a_{0010}^y y_{(n)} + a_{0020}^y y_{(n)}^2 + a_{0030}^y y_{(n)}^3 + a_{0040}^y y_{(n)}^4 + \dots \\ &+ a_{0100}^y \theta_{(n)} + a_{0200}^y \theta_{(n)}^2 + a_{0300}^y \theta_{(n)}^3 + a_{0400}^y \theta_{(n)}^4 + \dots \\ &+ a_{1000}^y x_{(n)} + a_{2000}^y x_{(n)}^2 + a_{3000}^y x_{(n)}^3 + a_{4000}^y x_{(n)}^4 + \dots \\ &+ a_{1111}^y x_{(n)} \theta_{(n)} y_{(n)} \phi_{(n)}, \end{aligned} \tag{6}$$

where  $n = 1, 2, \dots, N$ , and  $N$  is the number of coincidence events used in the analysis. The overdetermined set of Eqs. (6) represents a direct comparison of the reconstructed vertex position  $y_{Tg}^{\text{Optics}}$  to the measured value  $y_{Tg}^{\text{Measured}}$ . Initially a consistent polynomial expansion to fourth degree ( $i+j+k+l \leq 4$ ) was considered, which depends on 70 matrix elements  $a_{ijkl}^y$ . Using this ansatz in Eq. (4) defines a  $\chi^2$ -minimization function, which serves as an input to the simplex method. To be certain that the minimization did not converge to one of the local minima, the robustness of this method was examined by checking the convergence of the minimization algorithm for a large number of randomly chosen initial sets of parameters (see Fig. 6).

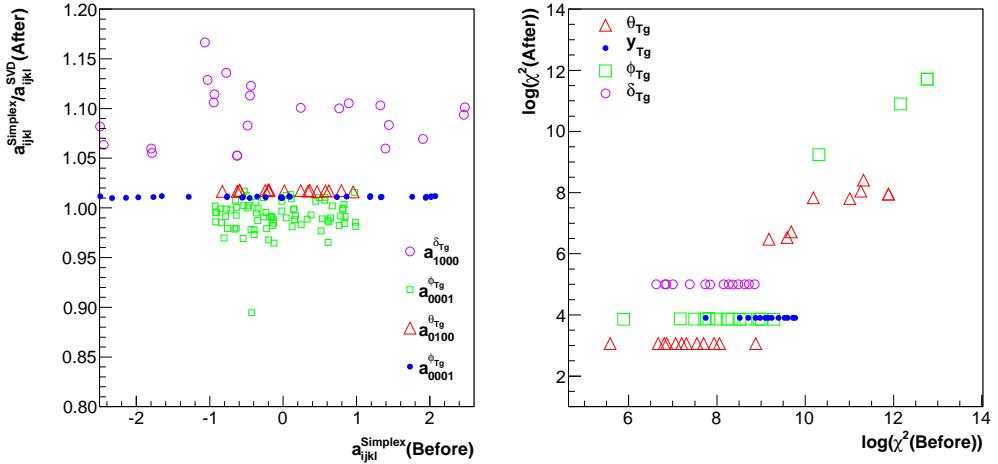


Figure 6: [Left] Robustness checks of the simplex minimization method for select matrix elements  $a_{ijkl}^{\Omega_{Tg}}$ . The analysis was done for a large set of randomly chosen initial conditions for each target coordinate. The fact that the vast majority of the initial conditions converge to a single value is an indication of the robustness of the method. [Right] The values of the  $\chi^2$ -function before and after simplex minimization for all four target coordinates. The method converges to a single  $\chi^2$  value for a wide range of initial conditions (note the log scales). The solution with the smallest  $\chi^2$  represents the result used in the optics-matrix.

The results were considered to be stable if the  $\chi^2$  defined by Eq. (4) converged to the same value for the majority of initial conditions. Small variations in  $\chi^2$  were allowed: they are caused by small matrix elements which are irrelevant for  $y_{Tg}$ , but have been set to non-zero values in order to additionally minimize  $\chi^2$  in a particular minimization process.

These matrix elements could be easily identified and excluded during the robustness checks because they are unstable and converge to a different value in each minimization. Ultimately only 25 matrix elements that had the smallest fluctuations were kept for the  $y_{Tg}$  matrix.

The SVD method was used next. To compute the matrix elements for  $y_{Tg}$ , the linear set of Eqs. (6) first needs to be rewritten in the form  $A \vec{a} = \vec{b}$  used in Eq. (5):

$$\begin{pmatrix} 1 & \phi_{(1)} & \cdots & x_{(1)}\theta_{(1)}y_{(1)}\phi_{(1)} \\ 1 & \phi_{(2)} & \cdots & x_{(2)}\theta_{(2)}y_{(2)}\phi_{(2)} \\ 1 & \phi_{(3)} & \cdots & x_{(3)}\theta_{(3)}y_{(3)}\phi_{(3)} \\ \vdots & \vdots & \ddots & \vdots \\ 1 & \phi_{(N-2)} & \cdots & x_{(N-2)}\theta_{(N-2)}y_{(N-2)}\phi_{(N-2)} \\ 1 & \phi_{(N-1)} & \cdots & x_{(N-1)}\theta_{(N-1)}y_{(N-1)}\phi_{(N-1)} \\ 1 & \phi_{(N)} & \cdots & x_{(N)}\theta_{(N)}y_{(N)}\phi_{(N)} \end{pmatrix} \begin{pmatrix} a_{0000} \\ a_{0001} \\ \vdots \\ a_{1111} \end{pmatrix} = \begin{pmatrix} y_{Tg(1)} \\ y_{Tg(2)} \\ y_{Tg(3)} \\ \vdots \\ y_{Tg(N-2)} \\ y_{Tg(N-1)} \\ y_{Tg(N)} \end{pmatrix},$$

where  $\vec{a}$  contains  $M$  unknown matrix elements  $a_{ijkl}^y$  to be determined by the SVD,  $\vec{b}$  contains  $N$  measured values of  $y_{Tg}$ , and  $A$  is filled with the products of detector variables accompanying the matrix elements in the polynomial expansion of Eq. (6) for each event.

The SVD analysis also began with 70 matrix elements, but was not applied to one combined data set as in the simplex method in order to extract the most relevant ones. Rather, it was used on each set of data separately. From the comparison of the matrix elements obtained with different calibration data sets, only the elements fluctuating by less than 100 % were selected. Although this choice appears to be arbitrary, the results do not change much by modifying this criterion, for example, by including elements with as much as  $\pm 1000$  % fluctuation. The final set of matrix elements contained only 37 of the best entries. With these elements, the entire analysis was repeated in order to calculate their final values. The most relevant elements are listed in Table 2. The result of the calibration of  $y_{Tg}$  is shown in Fig. 7.

### 5.2. Angular coordinates

For the calibration of the angular variables  $\theta_{Tg}$  and  $\phi_{Tg}$ , a set of quasi-elastic data on carbon and deuterium targets taken with the sieve-slit collimator was analyzed. The

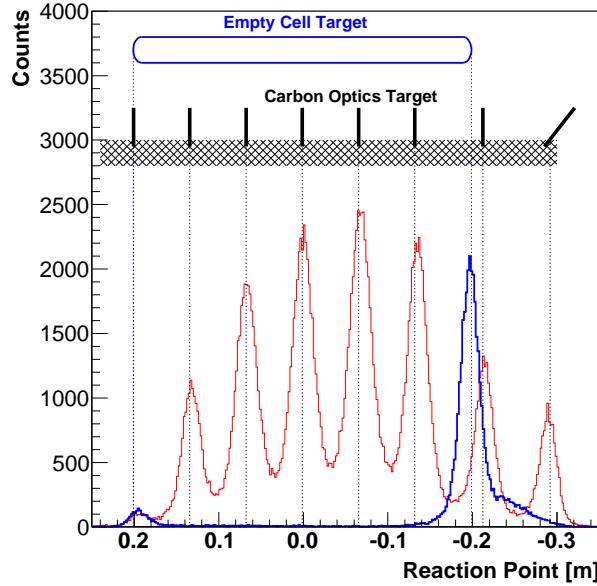


Figure 7: The reconstructed vertex position (reaction point) for the multi-foil carbon target and the empty cell of the production target, by using the SVD technique. The vertical dashed lines indicate the actual positions of the carbon foils and the empty-cell glass windows. The small shoulder to the right of the reconstructed empty-cell entry window is due to the jet of  ${}^4\text{He}$  gas used to cool the window at the beam impact point.

particles that pass through different holes can be well separated and localized at the detector plane.

By knowing the detector coordinates and the accurate position of the corresponding hole in the sieve, the target variables can be calculated. From the reaction point at the target (see Fig. 8),  $\theta_{\text{Tg}}$  and  $\phi_{\text{Tg}}$  can be calculated:

$$\tan \phi_{\text{Tg}} = \frac{y_{\text{Sieve}} - y_{\text{Tg}}}{z_{\text{Sieve}} - z_{\text{Tg}}} , \quad \tan \theta_{\text{Tg}} = \frac{x_{\text{Sieve}} - x_{\text{Tg}}}{z_{\text{Sieve}} - z_{\text{Tg}}} .$$

By using the values of the target variables, a set of linear equations has been written for all measured events, and matrix elements determined by using both numerical approaches. In the simplex method, 30 matrix elements for  $\theta_{\text{Tg}}$  and 68 elements for  $\phi_{\text{Tg}}$  were retained. Robustness checks for both angular variables were repeated to ensure that the global minimum had been reached.

The SVD analysis also started with 70 matrix elements, which were ultimately re-

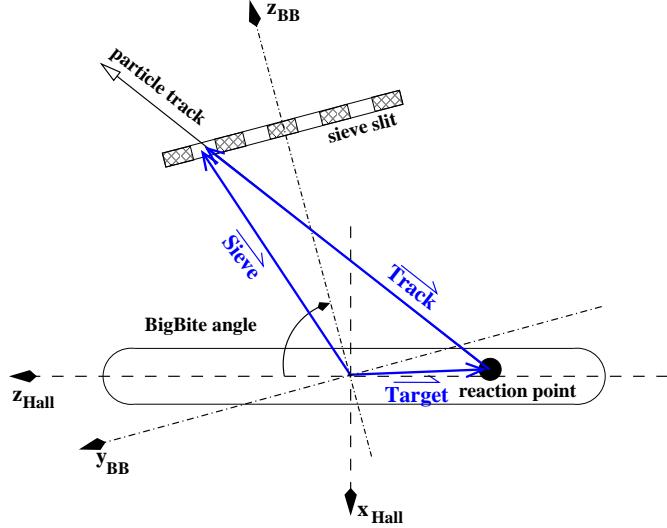


Figure 8: Position of the sieve-slit collimator relative to the target. The vector of the particle track through a particular hole in the sieve is the difference of the position vector at the hole and the reaction-point vector. BigBite is positioned at  $-75^\circ$  with respect to the beam direction. Other settings are listed in Table 1.

duced to 37 for  $\theta_{Tg}$  and 51 for  $\phi_{Tg}$ , again taking into account only those elements that fluctuated by less than 100 %. Figure 3 (right) shows the reconstructed sieve pattern. The majority of the holes are reconstructed, except those obscured by parts of the experimental apparatus due to specific geometric constraints during the experiment. In order to demonstrate the effect of gradually excluding redundant matrix elements, Fig. 9 shows the reconstructed top row of the sieve-slit collimator holes when the elements with up to  $\pm 1000\%$ ,  $\pm 100\%$ , and  $\pm 20\%$  fluctuations are retained. There is virtually no difference in the reconstructed pattern when all elements exceeding the  $\pm 100\%$  fluctuations are dropped, while errors start to appear when those fluctuating by less than  $\pm 100\%$  are dropped.

The quality of the sieve-pattern reconstruction was examined by comparing the centers of the reconstructed holes with their true positions. Figure 10 shows that, with the exception of a few holes near the acceptance edges, these deviations are smaller than 2 mm in the vertical, and smaller than 4 mm in the horizontal direction. This is much less than the hole diameter, which is 19.1 mm.

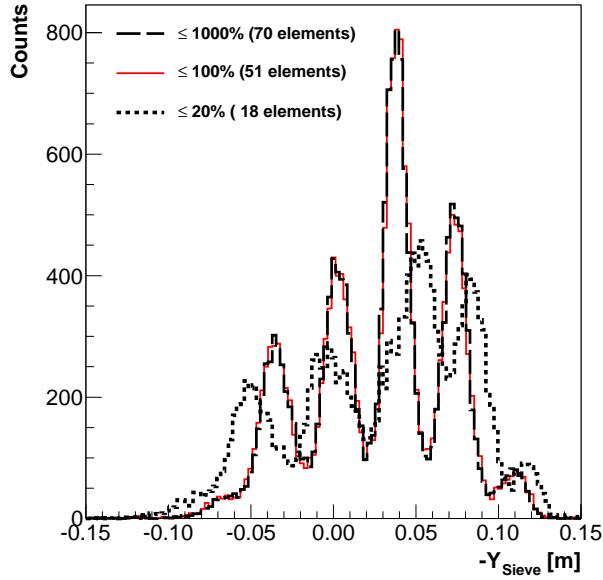


Figure 9: The reconstructed positions of the holes in the top row of the sieve-slit collimator, computed from  $\phi_{Tg}$ . The quality of the reconstruction depends on the number of included matrix elements. There is almost no difference when the elements fluctuating by up to  $\pm 1000\%$  are retained (70 elements, dashed lines) or only those that fluctuate by up to  $\pm 100\%$  (51 elements, full line). The quality deteriorates if too many elements are dropped (i.e. keeping 18 elements fluctuating by less than  $\pm 20\%$ , dotted lines).

Once the sieve pattern was reconstructed, an absolute calibration had to be performed to correct for any BigBite misalignment and mispointing. For that purpose hydrogen and deuterium elastic data were used. By comparing the direction of the momentum transfer vector from the HRS to the calculated values of  $\theta_{Tg}$  and  $\phi_{Tg}$ , the zero-order matrix elements could be properly determined and the offsets corrected. In addition, the precise distance between the target and the sieve-slit collimator was obtained, which we were not able to measure precisely due to physical obstacles between the target and BigBite. From this analysis, the sieve slit was determined to be positioned 1.13 m away from the target.

### 5.3. Momentum

The matrix elements for the  $\delta_{Tg}$  variable were obtained by using data from elastic scattering of electrons on hydrogen and deuterium for which the particle momentum in

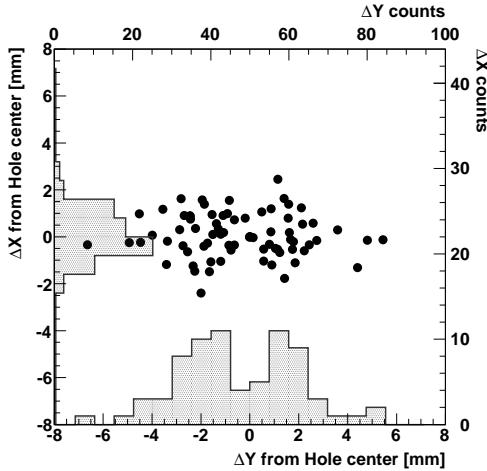


Figure 10: Distribution of vertical ( $\Delta X$ ) and horizontal ( $\Delta Y$ ) deviation of the center of each reconstructed sieve-slit hole from its true position. Observed deviations are much smaller than the diameter of a hole, which is 19.1 mm. The horizontal and vertical histograms (top and right axis labels, respectively) represent the distributions in the horizontal and vertical directions.

BigBite should be exactly the same as the momentum transfer  $\vec{q}$  given by the HRS. We assumed that  $\delta_{Tg}$  depends only on  $x_{\text{Det}}$  and  $\theta_{\text{Det}}$ , while the dependencies involving  $y_{\text{Det}}$  and  $\phi_{\text{Det}}$  were neglected. Furthermore, the use of in-plane coordinates in the analysis for  $\delta_{Tg}$  could result in an erroneous matrix due to the strong  $\phi_{Tg}$  dependence inherent to elastic scattering (events strongly concentrated at one edge of the acceptance). Considering only  $x_{\text{Det}}$  and  $\theta_{\text{Det}}$  matrix elements,  $\delta_{Tg}$  can be expressed as

$$\delta_{Tg} = \frac{q_{\text{HRS}} - \Delta_{\text{Loss}}}{p_c} - 1 = a_{0000}^\delta + a_{1000}^\delta x_{\text{Det}} + a_{0100}^\delta \theta_{\text{Det}} + \dots \quad (7)$$

In order to obtain the optics matrix applicable to all types of particles, energy losses  $\Delta_{\text{Loss}}$  for particle transport through the target enclosure and materials within the BigBite spectrometer were studied carefully. The energy losses were estimated by the Bethe-Bloch formula [29], but since the losses were significant, the formula had to be integrated over the complete particle track for each particle type and each initial momentum. The two largest contributions to the total momentum loss came from the target cell walls and from the air between the target and the detectors. (The latter losses could be alleviated by using a helium bag between the target and the detectors, but its benefits were considered

to be smaller than the technical problems involved.) The resulting corrections that were taken into account in Eq. (7) are shown in Fig. 11 (left).

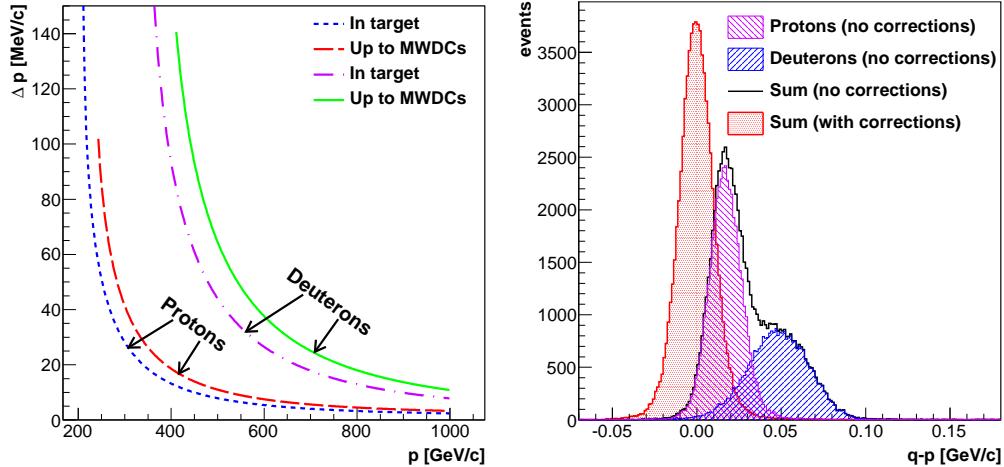


Figure 11: [Left] Momentum losses of protons and deuterons inside the target and the total momentum losses up to the MWDCs. [Right] Quality of reconstructed momentum for elastic protons and deuterons. If energy losses are not taken into account, two peaks are visible (center and right histograms summed to the full curve). With proper inclusion of energy losses both peaks merge into one (left histogram), resulting in better momentum resolution.

The elastic data available for calibration (momentum range approximately  $0.45 \text{ GeV}/c$  to  $0.7 \text{ GeV}/c$ ) covered only about half of the BigBite momentum acceptance. To calibrate the low-momentum region from  $0.2 \text{ GeV}/c$  to  $0.45 \text{ GeV}/c$ , we used protons from quasi-elastic scattering on  ${}^3\text{He}$  by exploiting the information from the scintillator dE- and E-planes; the deposited particle energy in each plane was directly mapped to the particle momentum, based on known properties of the scintillator material. The punch-through point, corresponding to the particular momentum at which the particle has just enough energy to penetrate through the scintillators, served as a reference.

Beside the proton punch-through point, two other points with exactly known energy deposits in the dE- and E-planes were identified, as illustrated in Fig. 12. With the additional information from these points, a complete momentum calibration was possible. To compute the  $\delta_{Tg}$  matrix elements, both numerical approaches described above were used. Since the available data were rather sparse, the search for the most stable matrix

elements was not performed and a complete expansion to fifth order was considered in both techniques. Since only a two-variable dependency was assumed, a complete description was achieved by using only 21 matrix elements.

The comparison of the most relevant matrix elements obtained from both numerical approaches is again shown in Table 2. Figure 12 (right) shows that the  $\delta_{Tg}$  matrix is well under control. The reconstructed momentum agrees with the simulation of energy losses inside the scintillation planes for the complete momentum acceptance of BigBite, for both protons and deuterons. Figure 5 shows the missing-mass peak for the  $^2\text{H}(\text{e}, \text{e}'\text{p})\text{n}$  process. The resolution of the reconstructed neutron mass is approximately  $4 \text{ MeV}/c^2$ .

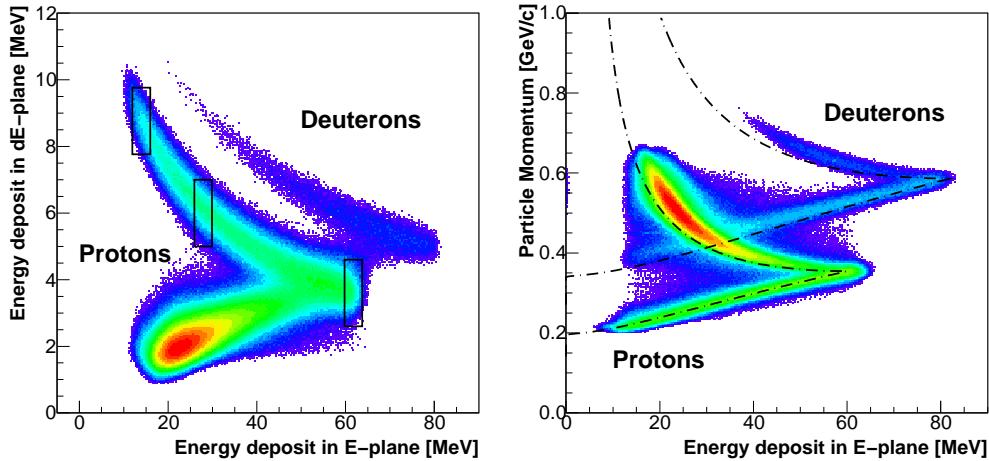


Figure 12: [Left] The energy losses in the thin (3 mm) scintillator dE-plane versus the energy losses in the thicker (3 cm) E-plane. The punch-through points, at which the protons and deuterons have just enough energy to penetrate both scintillation planes, are clearly visible. The black boxes show sections of events with precisely determined momenta that were used in the  $\delta_{Tg}$  calibration. [Right] Particle momentum as a function of energy losses in the E-plane for  $^3\text{He}$  data. The deuterons can be clearly distinguished from the protons. The measurements agree well with the simulation (dot-dashed line).

#### 5.4. Resolution

The quality of the BigBite optics was also studied. The resolution of the vertex position was estimated from the difference between the reconstructed  $y_{Tg}$  and the true position at the target by taking the width (sigma) of the obtained distribution. This part of the analysis was done by using 2-pass (2.425 GeV beam) quasi-elastic carbon

data. The extracted values for the resolution of  $y_{Tg}$  in different momentum bins can be parameterized as

$$\sigma_{y_{Tg}} \approx 0.01 \left( 1 + \frac{0.02}{p^4} \right),$$

where the particle momentum is in  $\text{GeV}/c$  and the result is in meters. It is best at the upper limit of the accepted momentum range (about  $p = 0.7 \text{ GeV}/c$ ) where it amounts to  $\sigma_{y_{Tg}} = 1.1 \text{ cm}$ . The deterioration of the resolution at lower momenta is due to multiple scattering [29] in the air between the scattering chamber and the MWDCs.

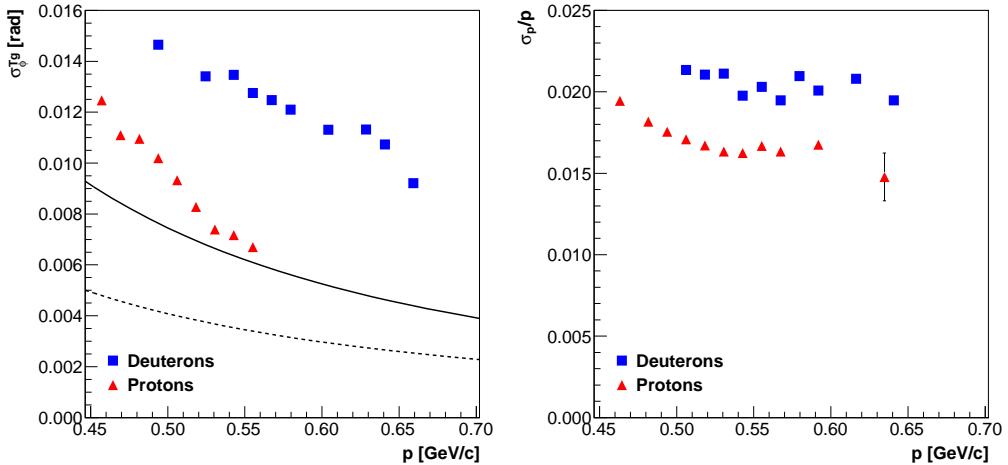


Figure 13: The absolute resolution of  $\phi_{Tg}$  and the relative momentum resolution as functions of the momentum measured by BigBite, obtained by the SVD method. Irreducible multiple-scattering contributions, mostly due to the air between the scattering chamber and MWDCs, are shown by full and dashed lines for deuterons and protons, respectively.

The resolutions of  $\theta_{Tg}$  and  $\phi_{Tg}$  were estimated by comparing them to the corresponding angles as determined from the momentum transfer  $\vec{q}$  in elastic scattering on hydrogen and deuterium. The direction of  $\vec{q}$  is given by the electron kinematics and determined by the HRS spectrometer. The corresponding HRS resolutions have been studied in [30]. Based on these values, the resolution of the reconstructed  $\vec{q}$  was estimated to be 6 mrad and 0.3 mrad for the vertical and horizontal angles, respectively. These contributions were subtracted in quadrature from the calculated peak widths, yielding the final resolutions attributable to BigBite. The result for  $\phi_{Tg}$  is shown in Fig. 13 (left). The strong momentum dependence of the resolution is again caused by multiple scattering in the tar-

get and the spectrometer. Different resolutions for deuterons and protons occur because the peak broadening in multiple scattering strongly depends on the particle mass (at a given momentum). As before, the biggest contributions come from the air. In a typical kinematics of the E05-102 experiment, the resolutions of  $\phi_{Tg}$  and  $\theta_{Tg}$  are  $\sigma_{\phi_{Tg}} \approx 7$  mrad and  $\sigma_{\theta_{Tg}} \approx 13$  mrad for  $0.55 \text{ GeV}/c$  protons, and approximately  $\sigma_{\phi_{Tg}} \approx 11$  mrad and  $\sigma_{\theta_{Tg}} \approx 13$  mrad for  $0.6 \text{ GeV}/c$  deuterons. (Due to multiple scattering, these resolutions are clearly much larger than the intrinsic MWDC resolutions mentioned in Section 2.)

The resolution of  $\delta_{Tg} = (p - p_c)/p_c$  was also determined from elastic data by comparing the magnitude of  $\vec{q}$  to the momentum reconstructed by BigBite. The analysis was done separately for the hydrogen and deuterium data sets. Figure 13 (right) shows the relative momentum resolution  $\sigma_p/p$  as a function of momentum. The relative momentum resolution is approximately 1.6 % for  $0.55 \text{ GeV}/c$  protons, and 2 % for  $0.6 \text{ GeV}/c$  deuterons. Figure 5 (right) shows the absolute resolution of  $\delta_{Tg}$ .

## 6. Summary

We have described the optics calibration of the BigBite spectrometer that was used to detect hadrons in the E05-102 experiment at Jefferson Lab. While the methods have been developed and applied to one spectrometer under very specific physical conditions, the same procedures can be applied to any spectrometer with a similar magnetic configuration and acceptance.

Two different approaches were considered: an analytical model that treats BigBite as an ideal dipole and a matrix formalism. The former approach results only in modest resolutions; still, resolutions of a few percent can be achieved by a suitable choice of parameters. The latter approach allows for a more precise calibration. Two numerical methods were used to determine the matrix elements, but the one based on singular value decomposition delivered better and more reliable results.

The vertex resolution for protons was found to be 1.2 cm at  $0.55 \text{ GeV}/c$  along the whole 40 cm target length. The resolution deteriorates significantly at lower momenta due to multiple scattering in the target, air, and detector material. The corresponding angular resolution is 7 mrad for the in-plane angle  $\phi_{Tg}$  and 13 mrad for the out-of-plane angle  $\theta_{Tg}$ . The angular resolution worsens at lower momenta due to multiple scattering,

with the effect more pronounced for deuterons. The relative momentum resolution for 0.55 GeV/ $c$  protons (best case) has been estimated to be 1.6%.

For 0.6 GeV/ $c$  deuterons (best case), we obtained the resolutions of 2% (momentum), 11 mrad ( $\phi_{Tg}$ ), and 13 mrad ( $\theta_{Tg}$ ).

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